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On an Integral Equation Arising in Inverse Scattering

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ON AN INTEGRAL EQUATION ARISING IN INVERSE SCATTERING

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ABSTRACT

A method of Fourier transforms and double series is applied for a solution of the following integral equation arising in inverse scattering:

$$\int_{-t}^x R(t+y)K(x,y)dy + K(x,t) + R(x+t) = 0, \quad t \leq x$$

where $R(w)$ is a Fourier transform of the given reflection coefficient in the differential equation $u''(k,x) + [k^2 - V(x)] u(k,x) = 0, \quad -\infty < x < \infty$.

By assuming $R(w)$ to be analytic, it is found that the unknown function can be constructed in the following forms:

$$i) \quad K(x,y) = \sum_{n=0}^{\infty} A_n(x) y^n, \quad ii) \quad K(x,y) = \sum_{p=0}^{\infty} \sum_{m+n=p} C_{m,n} x^m y^n$$

in which these series converge uniformly and absolutely for $|x| < x(R)$. $x(R)$ depends on $R(w)$. Therefore, if the reflection coefficient of the above differential equation is given with suitable conditions, $V(x)$ can be found by solving the integral equation and by the relation $V(x) = 2 \frac{d}{dx} K(x,x)$. Some related topics are discussed.

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1. Introduction

Let us consider the following differential equation,

$$\frac{d^2}{dx^2} U(k,x) + [k^2 - V(x)] U(k,x) = 0 ,$$

which appears in many physical problems.

If $V(x)$ is given, we can obtain, in principle, a solution such that

$$U(k,x) \sim e^{ikx} + r(k) e^{-ikx} , \quad x \rightarrow -\infty \quad (-\infty < x < \infty)$$

$$U(k,x) \sim t(k) e^{ikx} , \quad x \rightarrow +\infty$$

where $r(k)$, $t(k)$ are respectively reflection and transmission coefficients.*

Conversely, one can obtain the ionization density from a knowledge of the reflection coefficient. The latter data are obtainable from the time delay of a pulsed radio wave transmitted from the earth and reflected back by the ionosphere. This ionization density is calculated as follows:

a) we first determine the kernel $K(x,y)$ from the integral equation

$$(1) \quad \int_{-t}^x R(t+y) K(x,y) dy + K(x,t) + R(x+t) = 0, \quad t \leq x$$

where $R(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} r(k) e^{-ikt} dk$, $r(k)$ being a complex reflection coefficient.

* See [4] Appendix and [1] p. 5-7.

** See [1], section 3. We assume that $R(t)$ is analytic for $t \geq 0$ in this report unless specified.

b) We next calculate $V(x) = 2 \frac{d}{dx} K(x,x)$.

c) Finally $V(x)$ is proportional to the ionization density.

Equation (1) can be regarded as the following Fredholm integral equation

$$(2) \quad -\lambda \int_{-x}^x R(t+y) f(y) dy + f(t) + g(t) = 0 \quad (x = \text{fixed})$$

with the condition $R(s) = 0$ for $s < 0$; $\lambda = -1$, $f(y) = K(x,y)$ and $g(t) = R(x+t)$.

It follows that for every $t \in [-x, x]$, $f(t)$ has a solution belonging to L_2 (or C) if $R(s) \in L_2$ (or C), where L_2 is the class of square summable functions (C being the class of continuous functions).

Evaluation of a solution of (2) is very difficult, due to the condition $R(s) = 0$ for $s < 0$, by the standard Fredholm theory. It is therefore desirable to develop a new method of solving (2). From (1), after changing variables $z = t+y$, $w = x+t$, we obtain

$$(3) \quad \int_0^w R(z) K(x, x+z-w) dz + K(x, w-x) + R(w) = 0, \quad w \geq 0$$

which can also be written as follows:

$$(4) \quad \int_0^w R(w-u) K(x, x-u) du + K(x, w-x) + R(w) = 0, \quad w \geq 0.$$

Applying a Fourier transform to (4), for the intervals $(0, \infty)$ and $(-\infty, 0)$, we obtain the following equations, respectively:

$$(i) \quad r_+(a) \int_{-\infty}^x K(x,y) e^{-iay} dy + \int_x^{\infty} K(x,y) e^{iay} dy + e^{-iax} r_+(a) = 0$$

*
for $I_m(a) \gg 0$

*See Appendix, p. 26. $I_m(a)$ means the imaginary part of a ; $A \gg 0$ means A is positive and very large. Similarly $A \ll 0$ means A is negative and very large.

$$(ii) \quad -r_-(\alpha) \int_0^\infty K(x,y) e^{-i\alpha y} dy + \int_{-\infty}^{-x} K(x,y) e^{i\alpha y} dy + e^{-i\alpha x} r_-(\alpha) = 0$$

for $I_m(\alpha) \ll 0$

$$\text{where } r_+(\alpha) = \int_0^\infty R(t) e^{i\alpha t} dt = r(\alpha), \quad r_-(\alpha) = \int_{-\infty}^0 R(t) e^{i\alpha t} dt$$

Here, we have extended the definition of $R(t)$ for $t \leq 0$. **

Substituting $-\alpha$ for α into i) and ii), we obtain, respectively, the following:

$$(iii) \quad r_+(-\alpha) \int_{-\infty}^x K(x,y) e^{i\alpha y} dy + \int_{-x}^\infty K(x,y) e^{-i\alpha y} dy + e^{i\alpha x} r_+(-\alpha) = 0$$

for $I_m(\alpha) \ll 0$

$$(iv) \quad -r_-(-\alpha) \int_x^\infty K(x,y) e^{i\alpha y} dy + \int_{-\infty}^{-x} K(x,y) e^{-i\alpha y} dy + e^{i\alpha x} r_-(-\alpha) = 0$$

for $I_m(\alpha) \gg 0$.

The sets of equations (i,iv) and (ii,iii) cannot be solved uniquely unless we have more conditions on $K(x,y)$ since there are four unknowns but only two equations, respectively. Nevertheless, when $x = 0$, they are solvable for $K(0,y)$ and we obtain the following equations from (i-iv):

$$(i') \quad r_+(\alpha) \int_{-\infty}^0 K(0,y) e^{-i\alpha y} dy + \int_0^\infty K(0,y) e^{i\alpha y} dy + r_+(\alpha) = 0, \quad I_m(\alpha) \gg 0$$

$$(ii') \quad -r_-(\alpha) \int_0^\infty K(0,y) e^{-i\alpha y} dy + \int_{-\infty}^0 K(0,y) e^{i\alpha y} dy + r_-(\alpha) = 0, \quad I_m(\alpha) \ll 0$$

$$(iii') \quad r_+(-\alpha) \int_{-\infty}^0 K(0,y) e^{i\alpha y} dy + \int_0^\infty K(0,y) e^{-i\alpha y} dy + r_+(-\alpha) = 0, \quad I_m(\alpha) \ll 0$$

**

We take the analytical continuation of $R(t)$ as its extended definition for $t \leq 0$, in the above integral.

$$(iv') \quad -r_(-\alpha) \int_0^\infty K(0,y) e^{i\alpha y} dy + \int_{-\infty}^0 K(0,y) e^{-i\alpha y} dy + r_(-\alpha) = 0, \\ I_m(\alpha) \gg 0.$$

From (i') and (iv') we obtain

$$(5) \quad \hat{K}_+(0,\alpha) = \int_0^\infty K(0,y) e^{i\alpha y} dy = \begin{vmatrix} r_+(\alpha), r_+(\alpha) \\ 1, r_(-\alpha) \end{vmatrix} \begin{vmatrix} 1, r_+(\alpha) \\ -r_(-\alpha), 1 \end{vmatrix} \\ = r_+(\alpha) [r_(-\alpha) - 1] / [1 + r_+(\alpha) r_(-\alpha)] .$$

Similarly, from (iii') and (iiii'), we obtain

$$(6) \quad \hat{K}_-(0,\alpha) = \int_{-\infty}^0 K(0,y) e^{i\alpha y} dy = -r_(-\alpha) [1 + r_(-\alpha)] / [1 + r_(-\alpha) r_(-\alpha)] .$$

It follows that

$$(7) \quad K(0,y) = \frac{1}{2\pi} \int_{ic-\infty}^{ic+\infty} \hat{K}_+(0,\alpha) e^{-i\alpha y} d\alpha + \frac{1}{2\pi} \int_{id-\infty}^{id+\infty} \hat{K}_-(0,\alpha) e^{-i\alpha y} d\alpha \\ \text{for } c \gg 0 \text{ and } d \ll 0 .$$

According to Fredholm's theory, the solution of (2) can be written as follows:

$$f(t) = -g(t) - \int_{-x}^x K_R(t,y; \lambda) g(y) dy$$

where

$$K_R(t,y) = D(t,y; \lambda) / D(\lambda) \quad \text{if } D(\lambda) \neq 0 ,$$

$$D(\lambda) = 1 + \sum_{n=1}^{\infty} \frac{(-\lambda)^n}{n!} \int_{-x}^x \dots \int_{-x}^x R \begin{pmatrix} s_1, \dots, s_n \\ s_1, \dots, s_n \end{pmatrix} ds_1 \dots ds_n ,$$

$$D(t,y; \lambda) = 1 + \sum_{n=1}^{\infty} \frac{(-\lambda)^n}{n!} \int_{-x}^x \dots \int_{-x}^x R \begin{pmatrix} t, s_1, \dots, s_n \\ y, s_1, \dots, s_n \end{pmatrix} ds_1 \dots ds_n ,$$

and $R \begin{pmatrix} x_1, \dots, x_n \\ y_1, \dots, y_n \end{pmatrix} = \begin{vmatrix} R(x_1+y_1), \dots, R(x_1+y_n) \\ \vdots \\ R(x_n+y_1), \dots, R(x_n+y_n) \end{vmatrix}$

It follows that the solution of (1)

$$K(x,t) = -R(x+t) - \frac{1}{D(-1)} \int_{-t}^x D(t,y; -1) R(x+y) dy$$

is a function analytic with respect to t and meromorphic with respect to x if $R(w)$ is analytic with respect to w .

If $|\lambda_0(x)| > 1$ then $D(-1) \neq 0$, where $\lambda_0(x)$ is the first eigenvalue of (2), and this is always true if x is sufficiently small since $|\lambda_0(x)|^2$ is greater than $\left\{ \int_0^x \int_{-y}^x |R(t+y)|^2 dt dy \right\}^{-1}$.

It follows that $K(x,t)$ can always be expanded in a power series in x and t for small x . And $K(x,y)$ can also be expanded in a power series in x with coefficients as functions of y . In the following section, we shall discuss the two cases in which we could find the solution of (1) by a method of Fourier transforms.

2. The case where $K(x,y) = \sum_{n=0}^{\infty} A_n(y)x^n$

As in Section 1, we can assume that $K(x,y)$ can be expanded in power series of x : $K(x,y) = \sum_{n=0}^{\infty} A_n(y)x^n$ which is uniformly and absolutely convergent for small x , e.g. we can take any value of x

which is less than $x_0 = \sup \left\{ a \left| \int_0^a \int_{-y}^a |R(t+y)|^2 dt dy < 1 \right. \right\}$.

Theorem 1:

Let $R(t)$ be a given analytic function, then the integral equation (1) has a solution $K(x,y) = \sum_{n=0}^{\infty} A_n(y)x^n$ for $|x| < x_0$

where $A_n(y)$ can be obtained from the following recursion formula:

$$(8) \quad A_n(y) = \frac{1}{2\pi} \int_{ic-\infty}^{ic+\infty} \frac{r_+^0(\alpha)r_-^n(-\alpha) - r_+^n(\alpha)}{1 + r_+^0(\alpha)r_-^0(-\alpha)} e^{-iy\alpha} d\alpha \quad (c \gg 0)$$

$$- \frac{1}{2\pi} \int_{id-\infty}^{id+\infty} \frac{r_-^0(\alpha)r_+^n(-\alpha) + r_-^n(\alpha)}{1 + r_+^0(-\alpha)r_-^0(\alpha)} e^{-iy\alpha} d\alpha \quad (d \ll 0)$$

where

$$(9) \quad r_+^n(\alpha) = \int_0^{\infty} R^n(0,t) e^{i\alpha t} dt, \quad I_m(\alpha) \gg 0 \quad (n \geq 0)$$

$$(10) \quad r_-^n(\alpha) = \int_{-\infty}^0 R^n(0,t) e^{i\alpha t} dt, \quad I_m(\alpha) \ll 0$$

$$(11) \quad R^n(x,t) = \frac{1}{x} \left\{ \int_{-t}^x R(t+y) H_{n-1}(0,y) dy + H_{n-1}(0,t) + R^{n-1}(x,t) \right\}$$

$$(12) \quad H_n(x,t) = \frac{1}{x} \left\{ H_{n-1}(x,y) - H_{n-1}(0,y) \right\} \quad (n \geq 1)$$

$$(13) \quad R^0(x,y) = R(x+y)$$

$$(14) \quad H_0(x,y) = K(x,y).$$

It is known that when $n=0$, $A_0(y) = K(0,y)$ and formula (8) becomes

identical with (7); when $n=1$, $A_1(y) = H_1(0, y) = \left. \frac{\partial}{\partial x} K(x, y) \right|_{x=0}$, and

so on, and in general $A_n(y) = H_n(0, y) = \left. \left(\frac{\partial}{\partial x} \right)^n K(x, y) \right|_{x=0}$.

Proof of Theorem 1:

When $n=0$, we have already proven the theorem in Section 1.

By Mathematical Induction, suppose it has been proven for the cases $k < n$, then for $k=n$, we would have from (1)

$$(15) \quad \int_{-t}^x R(t+y) H_n(x, y) dy + H_n(x, t) + R^n(x, t) = 0.$$

This is true since, by our assumption, we already know that

$$(16) \quad \int_{-t}^x R(t+y) H_{n-1}(x, y) dy + H_{n-1}(x, t) + R^{n-1}(x, t) = 0$$

holds. It follows from (16) that

$$(17) \quad \begin{aligned} & \int_{-t}^x R(t+y) \left\{ H_{n-1}(x, y) - H_{n-1}(0, y) \right\} dy + H_{n-1}(x, t) - H_{n-1}(0, t) \\ & + R^{n-1}(x, t) + \int_{-t}^x R(t+y) H_{n-1}(0, y) dy + H_{n-1}(0, t) = 0. \end{aligned}$$

Dividing (17) by x , and using the formulas (11) and (12), we obtain the integral equation (15). Applying a method similar to that used in Section 1, we obtain

$$(I) \quad r_+^0(\alpha) \int_{-\infty}^x H_n(x, y) e^{-i\alpha y} dy + \int_{-x}^{\infty} H_n(x, y) e^{i\alpha y} dy + r_+^n(x, \alpha) = 0$$

$$\text{Im}(\alpha) \gg 0$$

$$(II) \quad -r_+^n(\alpha) \int_x^\infty H_n(x,y) e^{-i\alpha y} dy + \int_{-\infty}^{-x} H_n(x,y) e^{i\alpha y} dy + r_-^n(x,\alpha) = 0$$

$$\text{Im}(\alpha) \ll 0$$

where

$$r_+^n(x,\alpha) = \int_{-x}^\infty R^n(x,y) e^{i\alpha y} dy$$

$$r_-^n(x,\alpha) = \int_{-\infty}^{-x} R^n(x,y) e^{i\alpha y} dy .$$

Putting $x=0$, and using the same procedure as in Section 1, we obtain formula (8), where $r_+^n(\alpha) \equiv r_+^n(0,\alpha)$, $r_-^n(\alpha) \equiv r_-^n(0,\alpha)$.

3. The case where $K(x,y) = \sum_{n=0}^\infty \sum_{m+n=p} c_{m,n} x^m y^n$

Theorem 2:

Let $R(w) = \sum_{n=0}^\infty \gamma_n w^n$ be given, then the integral equation (1)

has a solution $K(x,y) = \sum_{p=0}^\infty \sum_{m+n=p} c_{m,n} x^m y^n$ for $|y| \leq |x| \leq x_0$, where

$c_{m,n}$ can be determined from γ_k by the following
 $(m,n=0,1,2,\dots)$ $(k=0,1,2,\dots)$

recursion formula:

$$(18) \quad \delta_{0,q} \gamma_l + \sum_{k=0}^{l-1} \sum_{m=0}^q (-1)^k \binom{k+q-m}{q-m} \frac{c_{m,k+q-m} \gamma_{l-k-1}}{l \binom{l-1}{k}} + \sum_{m=0}^q (-1)^{q-m} c_{m,l+q-m} \binom{l+q-m}{q-m} = 0$$

for $q,l = 0,1,2,3,\dots$ (Here, we use the convention $\sum_{k=0}^{-1} f_k = 0$, δ_{ij}

being Kronecker's delta).

The procedure to obtain $C_{m,n}$ from (18) is as follows:

Putting $q=0$, $l=0,1,2,3\dots$, we can obtain $C_{0,m}$ ($m=0,1,2,\dots$), step by step; then we obtain $C_{1,0}$ from $C_{0,0}$ and $C_{0,1}$, and so on. We obtain $C_{1,m}$ from $C_{0,p}$ ($0 \leq p \leq m+1$) and $C_{1,q}$ ($0 \leq q \leq m-1$). Similarly, we obtain $C_{m,n}$ from $C_{p,q}$ ($0 \leq p \leq m-1$, $0 \leq q \leq m+1$) and $C_{m,l}$ ($0 \leq l \leq n-1$).

Proof of Theorem 2:

Since $K(x,y) = \sum \sum C_{m,n} x^m y^n$, we obtain

$$\begin{aligned} K(x, x-u) &= \sum_{p=0}^{\infty} \sum_{m+n=p} C_{m,n} x^m (x-u)^n \\ &= \sum_{p=0}^{\infty} \sum_{m=0}^p C_{m,p-m} \sum_{q=m}^p (-1)^{p-q} \binom{p-m}{q-m} u^{p-q} x^q \\ &= \sum_{p=0}^{\infty} \sum_{q=0}^p \left(\sum_{m=0}^q C_{m,p-m} (-1)^{p-q} \binom{p-m}{q-m} u^{p-q} \right) x^q \\ &= \sum_{q=0}^{\infty} \left(\sum_{p=q}^{\infty} \sum_{m=0}^q C_{m,p-m} (-1)^{p-q} \binom{p-m}{q-m} u^{p-q} \right) x^q. \end{aligned}$$

Similarly

$$\begin{aligned} (19) \quad K(x, w-x) &= \sum_{p=0}^{\infty} \sum_{m+n=p} C_{m,n} x^m (w-x)^n \\ &= \sum_{q=0}^{\infty} \left(\sum_{p=q}^{\infty} \sum_{m=0}^q C_{m,p-m} (-1)^{q-m} \binom{p-m}{q-m} w^{p-q} \right) x^q, \end{aligned}$$

and

$$\begin{aligned} (20) \quad \int_0^w R(w-u) K(x, x-u) du &= \int_0^w R(w-u) \left\{ \sum_{q=0}^{\infty} \left(\sum_{p=q}^{\infty} \sum_{m=0}^q C_{m,p-m} (-1)^{p-q} \binom{p-m}{q-m} u^{p-q} \right) x^q \right\} du \\ &= \sum_{q=0}^{\infty} \left\{ \int_0^w R(w-u) \left[\sum_{p=q}^{\infty} \sum_{m=0}^q C_{m,p-m} (-1)^{p-q} \binom{p-m}{q-m} u^{p-q} \right] du \right\} x^q. \end{aligned}$$

Substituting (19) and (20) into (4) and arranging in powers of x , we obtain from the coefficient of q^{th} power the following:

$$\begin{aligned}
 & \int_0^w R(w-u) \left\{ \sum_{p=q}^{\infty} \sum_{m=0}^q c_{m,p-m} (-1)^{p-q} \binom{p-m}{q-m} u^{p-q} \right\} du \\
 & + \sum_{p=q}^{\infty} \sum_{m=0}^q c_{m,p-m} (-1)^{q-m} \binom{p-m}{q-m} w^{p-q} + \delta_{0,q} R(w) = 0.
 \end{aligned}
 \tag{21}$$

Also, we have

$$\begin{aligned}
 & \int_0^w R(w-u) \left\{ \sum_{p=q}^{\infty} \sum_{m=0}^q c_{m,p-m} (-1)^{p-q} \binom{p-m}{q-m} u^{p-q} \right\} du \\
 & = \sum_{p=q}^{\infty} \sum_{m=0}^q c_{m,p-m} (-1)^{p-q} \binom{p-m}{q-m} \int_0^w u^{p-q} \sum_{n=0}^{\infty} \gamma_n (w-u)^n du \\
 & \quad (\text{put } t=n+p-q) \quad = \sum_{p=q}^{\infty} \sum_{m=0}^q c_{m,p-m} (-1)^{p-q} \binom{p-m}{q-m} \sum_{n=0}^{\infty} \gamma_n \sum_{s=0}^n \binom{n}{s} (-1)^{n-s} \int_0^w u^{n-s+p-q} du \cdot w^s \\
 & \quad = \sum_{p=q}^{\infty} \sum_{m=0}^q c_{m,p-m} (-1)^{p-q} \binom{p-m}{q-m} \sum_{\ell=p-q}^{\infty} \gamma_{\ell+q-p} \sum_{s=0}^{\ell+q-p} \binom{\ell+q-p}{s} \frac{(-1)^{\ell+q+p+s} w^{\ell+1}}{w^{\ell-s-1}}
 \end{aligned}$$

$$\begin{aligned}
 & \quad (\text{put } k=p-q) \quad = \sum_{\ell=0}^{\infty} \left\{ \sum_{k=0}^{\ell} \sum_{m=0}^q c_{m,k+q-m} (-1)^k \binom{k+q-m}{q-m} \cdot \frac{\gamma_{\ell-k}}{(\ell+1) \binom{\ell}{k}} \right\} w^{\ell+1}.
 \end{aligned}$$

$$(\text{Here we use } (*) \dots \sum_{s=0}^{\ell-k} \binom{\ell-k}{s} \frac{(-1)^{\ell+k+s}}{w^{\ell-s+1}} = \frac{1}{(\ell+1) \binom{\ell}{k}} \dots)^*.$$

* See Appendix

Substituting the above expression into (21) and equating the coefficient of the t^{th} power of w , we obtain (18).

4. Remarks, Special Cases, and Examples.

A recursion formula equivalent to the identity (18) can be derived from equation (i) if $K(x,y)$ satisfies some conditions such that integrations by parts are possible in equation (i). We have the following theorem:

Theorem 3: If $\left(\frac{\partial}{\partial y}\right)^n K(x,y) = O(e^{cy})$ for a fixed constant c and for all natural numbers n , then by integrations by parts we can obtain the following double recursion formula from equation (i):

$$(22) \quad \delta_{0,s} k_r - r! \sum_{n=r}^{s+r} (-1)^n \binom{n}{r} c_{s+r-n,n} - \sum_{m=0}^{r-1} \sum_{n=r-m-1}^{s+r-m-1} (r-m-1)! k_m \binom{n}{r-m-1} \cdot c_{s+r-m-1-n,n} = 0$$

where

$$r_+(\alpha) = \sum_{n=0}^{\infty} \frac{k_n}{(i\alpha)^n}, \quad K(x,y) = \sum_{p=0}^{\infty} \sum_{m+n=p} c_{m,n} x^m y^n \quad \text{being assumed.}$$

Proof: By integrating by parts N times, we obtain

$$\int_{-\infty}^x K(x,y) e^{-iay} dy = -e^{-i\alpha x} \left[\sum_{n=0}^{N-1} \left(\frac{\partial}{\partial y} \right)^n K(x,y) \right]_{y=x}^{y=x} (i\alpha)^{-n-1} + O(\alpha^{-N-1})$$

$$\int_{-x}^{\infty} K(x,y) e^{iay} dy = e^{-i\alpha x} \left[\sum_{n=0}^{N-1} \left(\frac{\partial}{\partial y} \right)^n K(x,y) \right]_{y=-x}^{y=x} (-i\alpha)^{-n-1} + O(\alpha^{-N-1}).$$

Equating the r^{th} power of $(i\alpha)^{-1}$ in the equation (i), we obtain

$$(23) \quad \left[k_r + (-1)^{r-1} \left(\frac{\partial}{\partial y} \right)^r K(x, y) \right]_{y=-x} - \sum_{m=0}^{r-1} k_m \left(\frac{\partial}{\partial y} \right)^{r-m-1} K(x, y) \Big|_{y=x} = 0 \quad (r \leq N)$$

Putting $K(x, y) = \sum_{p=0}^{\infty} \sum_{m+n=p} c_{m,n} x^m y^n$ and equating the s^{th} power of x

in the identity (23), we obtain the identity (22). If we assume

$$R(w) = \sum_{n=0}^{\infty} \gamma_n w^n \text{ then } r_+(\alpha) = \sum_{n=0}^{\infty} \frac{k_n}{(\alpha)^{n+1}} = \int_0^{\infty} R(w) e^{i\alpha w} dw$$

$$= \sum_{n=0}^{\infty} \frac{(-1)^{n+1} n! \gamma_n}{(\alpha)^{n+1}}$$

i.e.

$$(24) \quad k_n = (-1)^{n+1} n! \gamma_n .$$

Substituting the identity (24) for k_n into the identity (22) and dividing by $r!$ in the latter, we can easily obtain the identity (18) if we make a suitable change of variables.

In Theorem 2 (or Theorem 3), computation of $c_{m,n}$ is sometimes tedious; e.g. $c_{m,n}$ ($m+n \leq 5$) in terms of γ_i are as follows:

[We could obtain $c_{m,n}$ in terms of k_i if we used the identity (24).]

$$c_{0,0} = -\gamma_0$$

$$c_{0,1} = -\gamma_1 + \gamma_0^2$$

$$c_{0,2} = -\gamma_2 + \frac{1}{2}\gamma_0^3$$

$$c_{0,3} = -\gamma_3 + \frac{2}{3}\gamma_2\gamma_0 - \frac{1}{6}\gamma_1^2 + \frac{1}{6}\gamma_1\gamma_0^2 - \frac{1}{6}\gamma_0^4$$

$$c_{0,4} = -\gamma_4 + \frac{1}{4}\gamma_2\gamma_0^2 - \frac{1}{24}\gamma_1^2\gamma_0 - \frac{1}{24}\gamma_0^5$$

$$c_{0,5} = -\gamma_5 + \frac{2}{5}\gamma_4\gamma_0 - \frac{1}{10}\gamma_3\gamma_1 + \frac{1}{20}\gamma_3\gamma_0^2 + \frac{1}{30}\gamma_2^2 + \frac{1}{30}\gamma_2\gamma_1\gamma_0 - \frac{1}{15}\gamma_2\gamma_0^3 - \frac{1}{120}\gamma_1^3 + \frac{1}{60}\gamma_1^2\gamma_0^2 - \frac{1}{120}\gamma_1\gamma_0^4 + \frac{1}{120}\gamma_0^6$$

$$c_{1,o} = c_{o,1} = -\gamma_1 + \gamma_o^2$$

$$c_{1,1} = -2\gamma_2 + 2\gamma_1\gamma_o - \gamma_o^3$$

$$c_{1,2} = -3\gamma_3 + \frac{1}{2}\gamma_1^2 + \frac{1}{2}\gamma_1\gamma_o^2 - \frac{1}{2}\gamma_o^4$$

$$c_{1,3} = -4\gamma_4 + 2\gamma_3\gamma_o - \frac{1}{3}\gamma_2\gamma_o^2 + \frac{1}{6}\gamma_1^2\gamma_o - \frac{1}{3}\gamma_1\gamma_o^3 + \frac{1}{6}\gamma_o^5$$

$$c_{1,4} = -5\gamma_5 + \frac{1}{2}\gamma_3\gamma_1 + \frac{1}{4}\gamma_3\gamma_o^2 - \frac{1}{6}\gamma_2^2 + \frac{1}{6}\gamma_2\gamma_1\gamma_o - \frac{1}{6}\gamma_2\gamma_o^2 - \frac{1}{24}\gamma_1^3$$

$$- \frac{1}{24}\gamma_1\gamma_o^4 + \frac{1}{24}\gamma_o^6$$

$$c_{2,o} = -\gamma_2 + 2\gamma_1\gamma_o - \frac{3}{2}\gamma_o^3$$

$$c_{2,1} = -3\gamma_3 + 2\gamma_2\gamma_o + \frac{3}{2}\gamma_1^2 - \frac{7}{2}\gamma_1\gamma_o^2 + \frac{3}{2}\gamma_o^4$$

$$c_{2,2} = -6\gamma_4 + 2\gamma_2\gamma_1 - \frac{1}{2}\gamma_2\gamma_o^2 - \frac{1}{4}\gamma_1^2\gamma_o - \gamma_1\gamma_o^3 + \frac{3}{4}\gamma_o^5$$

$$c_{2,3} = -10\gamma_5 + 4\gamma_4\gamma_o + \gamma_3\gamma_1 - \frac{3}{2}\gamma_3\gamma_o^2 + \frac{1}{3}\gamma_2^2 - \gamma_2\gamma_1\gamma_o + \frac{2}{3}\gamma_2\gamma_o^3$$

$$+ \frac{1}{4}\gamma_1^3 - \frac{1}{2}\gamma_1^2\gamma_o^2 + \frac{7}{12}\gamma_1\gamma_o^4 - \frac{1}{4}\gamma_o^6$$

$$c_{3,o} = -\gamma_3 + \frac{8}{3}\gamma_2\gamma_o + \frac{5}{6}\gamma_1^2 - \frac{23}{6}\gamma_1\gamma_o^2 + \frac{11}{6}\gamma_o^4$$

$$c_{3,1} = -4\gamma_4 + 2\gamma_3\gamma_o + 4\gamma_2\gamma_1 - \frac{13}{3}\gamma_2\gamma_o^2 - \frac{23}{6}\gamma_1^2\gamma_o + \frac{17}{3}\gamma_1\gamma_o^3 - \frac{11}{6}\gamma_o^5$$

$$c_{3,2} = -10\gamma_5 + 3\gamma_3\gamma_1 - \frac{3}{2}\gamma_3\gamma_o^2 + \frac{7}{3}\gamma_2^2 - \frac{7}{3}\gamma_2\gamma_1\gamma_o - \frac{1}{3}\gamma_2\gamma_o^3 - \frac{5}{12}\gamma_1^3$$

$$+ \frac{23}{12}\gamma_1\gamma_o^4 - \frac{11}{12}\gamma_o^6$$

$$c_{4,o} = -\gamma_4 + 4\gamma_3\gamma_o + 2\gamma_2\gamma_1 - \frac{53}{12}\gamma_2\gamma_o^2 - \frac{27}{8}\gamma_1^2\gamma_o + \frac{19}{3}\gamma_1\gamma_o^3 - \frac{19}{8}\gamma_o^5$$

$$c_{4,1} = -5\gamma_5 + 2\gamma_4\gamma_o + \frac{11}{2}\gamma_3\gamma_1 - \frac{23}{4}\gamma_3\gamma_o^2 + \frac{17}{6}\gamma_2^2 - \frac{55}{6}\gamma_2\gamma_1\gamma_o$$

$$+ \frac{19}{3}\gamma_2\gamma_o^3 - \frac{41}{24}\gamma_1^3 + \frac{35}{4}\gamma_1^2\gamma_o - \frac{209}{24}\gamma_1\gamma_o^4 + \frac{19}{8}\gamma_o^6$$

$$c_{5,o} = -\gamma_5 + \frac{32}{5}\gamma_4\gamma_o + \frac{29}{10}\gamma_3\gamma_1 - \frac{119}{20}\gamma_3\gamma_o^2 + \frac{31}{30}\gamma_2^2 - \frac{239}{30}\gamma_2\gamma_1\gamma_o$$

$$+ \frac{223}{30}\gamma_2\gamma_o^3 - \frac{121}{120}\gamma_1^3 + \frac{124}{15}\gamma_1^2\gamma_o - \frac{1201}{120}\gamma_1\gamma_o^4 + \frac{361}{120}\gamma_o^6.$$

$c_{q,l}$ can also be written in determinant form as follows:

$$(25) \quad c_{q,l} = (-1)^{l-1} \begin{vmatrix} q^f_0 & 1 & 0 & 0 & \dots & \dots & \dots & \dots & 0 \\ q^f_1 & \gamma_0 & 1 & 0 & \dots & \dots & \dots & \dots & \dots \\ q^f_2 & \frac{\gamma_1}{2} & \frac{-\gamma_0}{2} & 1 & \dots & \dots & \dots & \dots & \dots \\ \vdots & \vdots & \vdots & \vdots & \ddots & \ddots & \ddots & \ddots & \ddots \\ \vdots & \vdots & \vdots & \vdots & \ddots & \ddots & \ddots & \ddots & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ \vdots & \vdots & \vdots & \vdots & \ddots & \ddots & \ddots & \ddots & 1 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ q^f_l & \frac{\gamma_l}{l} & \frac{-\gamma_{l-1}}{l \binom{l}{0}} & \frac{\gamma_{l-2}}{l \binom{l-1}{1}} & \dots & \dots & \dots & \dots & \frac{(-1)^{l-1} \gamma_0}{l \binom{l-1}{l-1}} \end{vmatrix}$$

where

$$q^f_l = \delta_{0,q} \gamma_l + (-1)^q \sum_{m=0}^{q-1} (-1)^m \binom{l+q-m}{q-m} c_{m,l+q-m} + \sum_{k=0}^{l-1} \frac{(-1)^k \gamma_{l-k-1}}{l \binom{l-1}{k}} \cdot \sum_{m=0}^{q-1} \binom{k+q-m}{q-m} c_{m,k+q-m} \quad (q, l = 0, 1, 2, 3, \dots).$$

The formula (25) can be derived from the following formula (26) which is a different expression of (18), by applying Cramer's Rule for linear equations and by regarding $c_{q,k}$ ($0 \leq k \leq l$) as unknowns.

$$(26) \quad c_{q,t} + \sum_{k=0}^{t-1} (-1)^k \frac{\gamma_{t-k-1}}{t \binom{t-1}{k}} c_{q,k} + q^f_t = 0 \quad (0 \leq t \leq l).$$

It is known that we can obtain $C_{o,m}$ easily by calculating (25); nevertheless it is convenient to calculate $C_{p,q}$ for $p \geq 1$ by using (18).

The integral equations (3) and (4) are of Volterra type if we regard $K(x,y)$ as a given kernel and $R(x)$ as unknown; therefore, the solution of (3) and (4) [as well as (1)] can be said to be an Inverse Problem of the Volterra equation. For a given $K(x,w-x)$ in (3) [and in (4)], solving for $R(w)$, we obtain a solution $R(w)$ which in general not only depends on w but also depends on x . Consequently, the condition that $R(w)$ is independent of x is rather strong. It implies that functions $K(x,y)$ have a particular form. We shall discuss this situation in the following special cases:

(a) When $K(x,y) = A(x)B(y)$, we have from the identities (23) and (24), the following:

$$(a,1) \quad \gamma_0 = -k_0 = -K(x,-x) = -A(x)B(-x)$$

$$(a,2) \quad \gamma_1 = k_1 = -\gamma_0 A(x)B(x) - A(x)B'(x)$$

$$(a,3) \quad \gamma_2 = \frac{-1}{2} k_2 = -\frac{1}{2} \gamma_1 A(x)B(x) + \frac{1}{2} \gamma_0 A(x)B'(x) - \frac{1}{2} A(x)B''(-x).$$

From (a,1) and (a,2) we obtain, by eliminating $A(x)$,

$$(a,4) \quad \gamma_1 B(-x) = \gamma_0^2 B(x) + \gamma_0 B'(x) .$$

In changing the sign of x , we have

$$(a,5) \quad \gamma_1 B(x) = \gamma_0^2 B(-x) + \gamma_0 B'(x) .$$

By adding (a,4) and (a,5) and by subtracting (a,4) from (a,5) we obtain, respectively:

$$(a,6) \quad \gamma_0 \left\{ B'(x) + B'(-x) \right\} = -(\gamma_0^2 - \gamma_1^2) \left\{ B(x) + B(-x) \right\}$$

$$(a,7) \quad \gamma_0 \left\{ B'(x) - B'(-x) \right\} = (\gamma_0^2 + \gamma_1^2) \left\{ B(x) - B(-x) \right\} .$$

$$\text{By putting } G(x) = \frac{1}{2} \left\{ B(x) + B(-x) \right\}, \quad F(x) = \frac{1}{2} \left\{ B(x) - B(-x) \right\} ,$$

and differentiating (a,6) and (a,7), we obtain

$$(a,8) \quad \gamma_0^2 F''(x) + (\gamma_0^4 - \gamma_1^2) F(x) = 0$$

$$(a,9) \quad \gamma_0^2 G''(x) + (\gamma_0^4 - \gamma_1^2) G(x) = 0 .$$

By noticing that $G(x)$ is even and $F(x)$ is odd, and using (a,6) and

(a,7), we obtain ($C = \text{constant}$)

$$G(x) = \begin{cases} \frac{a}{\gamma_0} x + \frac{-a}{\gamma_0} x \\ C(e^{\frac{a}{\gamma_0} x} + e^{\frac{-a}{\gamma_0} x}) \end{cases} \text{ if } |\gamma_1| > \gamma_0^2 \text{ and } a = \sqrt{\gamma_1^2 - \gamma_0^4}$$

$$\begin{cases} \frac{a}{\gamma_0} ix + \frac{-a}{\gamma_0} ix \\ C(e^{\frac{a}{\gamma_0} ix} + e^{\frac{-a}{\gamma_0} ix}) \end{cases} \text{ if } |\gamma_1| < \gamma_0^2 \text{ and } a = \sqrt{\gamma_0^4 - \gamma_1^2}$$

$$F(x) = \begin{cases} \frac{\gamma_1 - \gamma_0^2}{a} C(e^{\frac{a}{\gamma_0} x} - e^{\frac{-a}{\gamma_0} x}) \text{ if } |\gamma_1| > \gamma_0^2 \text{ and } a = \sqrt{\gamma_1^2 - \gamma_0^4} \\ -\frac{\gamma_0^2 - \gamma_1}{ai} C(e^{\frac{a}{\gamma_0} ix} - e^{\frac{-a}{\gamma_0} ix}) \text{ if } |\gamma_1| < \gamma_0^2 \text{ and } a = \sqrt{\gamma_0^4 - \gamma_1^2} . \end{cases}$$

Consequently, we obtain

$$(26) \quad K(x, y) = \frac{-\gamma_0 B(y)}{B(-x)} = \begin{cases} \frac{(a + \gamma_1 - \gamma_0^2)e^{-\frac{a}{\gamma_0} y} + (a - \gamma_1 + \gamma_0^2)e^{\frac{a}{\gamma_0} y}}{\gamma_0 - \frac{a}{\gamma_0} x} \text{ if } \begin{cases} |\gamma_1| > \gamma_0^2 \\ a = \sqrt{\gamma_1^2 - \gamma_0^4} \end{cases} \\ \frac{(a - \gamma_1 + \gamma_0^2)e^{-\frac{a}{\gamma_0} y} + (a + \gamma_1 - \gamma_0^2)e^{\frac{a}{\gamma_0} y}}{\gamma_0 - \frac{a}{\gamma_0} x} \\ -\gamma_0 \frac{(\gamma_0^2 - \gamma_1) \sin(\frac{a}{\gamma_0} y) - a \cos(\frac{a}{\gamma_0} y)}{(\gamma_0^2 - \gamma_1) \sin(\frac{a}{\gamma_0} x) + a \cos(\frac{a}{\gamma_0} x)} \text{ if } \begin{cases} \gamma_0^2 > |\gamma_1| \\ a = \sqrt{\gamma_0^4 - \gamma_1^2} \end{cases} \\ -\gamma_0 \quad \text{if } \gamma_1 = \gamma_0^2 \\ -\gamma_0 \frac{2\gamma_0 y - 1}{2\gamma_0 x + 1} \text{ if } \gamma = -\gamma_0^2 \end{cases}$$

since for $\gamma_1 = \gamma_0^2$, (a,6), (a,7) become, respectively,

$$F'(x) = 0 \quad \text{and} \quad G'(x) = \gamma_0 F(x) .$$

It follows that $F(x) = C = 0$ because $F(x)$ is odd and hence $G(x) = 0$ - similarly for the case $\gamma_1 = -\gamma_0^2$. If $\gamma_0 = 0$, there is only the trivial solution $K(x,y) \equiv 0$. By the relation (a,2), if $\gamma_1 = 0$, we have $-\gamma_0 B(x) + B'(-x) = 0$ from which we can show that $\gamma_n = 0$ for $n > 1$. From (a,3) we find that $\gamma_2 = \gamma_1^2/\gamma_0$ and, in general, can show that

$$\gamma_n = \gamma_0 \left(\frac{\gamma_1}{\gamma_0} \right)^n, \text{ i.e., } R(x) = \gamma_0 e^{\frac{\gamma_1}{\gamma_0} x} \quad \text{Conversely, if } R(x) = \gamma_0 e^{\frac{\gamma_1}{\gamma_0} x} \text{ then}$$

$K(x,y)$ has exactly the form of (26); we shall show this in Example I.

(b) When $K(x,y) = A(x)B(y) + C(x)$, we have from the identities (23) and (24), the following:

$$(b,1) \quad -\gamma_0 = A(x)B(-x) + C(x)$$

$$(b,2) \quad -\gamma_1 = A(x)B'(-x) + \gamma_0 \left\{ A(x)B(x) + C(x) \right\}$$

$$(b,3) \quad -2\gamma_2 = A(x)B''(-x) - \gamma_0 A(x)B'(x) + \gamma_1 \left\{ A(x)B(x) + C(x) \right\} .$$

Eliminating $A(x)$, $C(x)$ from (b,1-3), we obtain

$$\begin{vmatrix} B(-x) & , & 1, & \gamma_0 \\ B'(-x) + \gamma_0 B(x) & , & \gamma_0, & \gamma_1 \\ B''(-x) - \gamma_0 B'(x) + \gamma_1 B(x), & \gamma_1, & 2\gamma_2 \end{vmatrix} = 0$$

which can be written as

$$(b,4) \quad B''(-x) - \gamma_0 B'(x) - \frac{2\gamma_2 - \gamma_0 \gamma_1}{\gamma_1 - \gamma_0} B'(-x) - \frac{2\gamma_0 \gamma_2 - \gamma_1^2}{\gamma_1 - \gamma_0} \left\{ B(x) - B(-x) \right\} = 0$$

By changing the sign of x and putting $K = \frac{2\gamma_2 - \gamma_0 \gamma_1}{\gamma_1 - \gamma_0}$, we have

$$B''(x) - \gamma_0 B'(-x) - KB'(x) + (\gamma_1 - \gamma_0 K) \left\{ B(-x) - B(x) \right\} = 0$$

By a procedure similar to the one used in (a), we obtain

$$(b,5) \quad G''(x) - (\gamma_0 + K) F'(x) = 0$$

$$(b,6) \quad F''(x) + (\gamma_0 - K) G'(x) + 2(\gamma_0 K - \gamma_1) F(x) = 0$$

$$\text{where } G(x) = \frac{1}{2} \left\{ B(x) + B(-x) \right\}, \quad F(x) = \frac{1}{2} \left\{ B(x) - B(-x) \right\}.$$

Differentiating (b,5-6) and eliminating $G''(x)$, we have

$$H''(x) - \lambda^2 H(x) = 0$$

where

$$H(x) = F'(x), \quad -\lambda^2 = \gamma_0^2 - 2\gamma_1 + 2K\gamma_0 - K^2.$$

It follows that if $\lambda^2 \neq 0$, we have a solution

$$B(x) = C_1 e^{\lambda x} + C_2 e^{-\lambda x} \quad (C_1, C_2 = \text{constants}).$$

And $A(x)$, $C(x)$ can be computed from $B(x)$ by solving simultaneous linear equations (b, 1-2) with respect to $A(x)$ and $C(x)$, i.e.,

$$A(x) = \frac{\begin{vmatrix} \gamma_0 & 1 \\ \gamma_1 & \gamma_0 \end{vmatrix}}{\begin{vmatrix} -B(-x), 1 \\ -B'(x) + \gamma_0 B(x), \gamma_0 \end{vmatrix}} = \frac{\gamma_0^2 - \gamma_1^2}{B'(-x) - \gamma_0 \{B(x) + B(-x)\}}$$

$$C(x) = \frac{\begin{vmatrix} -B(-x), \gamma_0 \\ -B'(-x) + \gamma_0 B(x), \gamma_1 \end{vmatrix}}{\begin{vmatrix} -B(-x), 1 \\ -B'(-x) + \gamma_0 B(x), \gamma_0 \end{vmatrix}} = \frac{-\gamma_0 B'(-x) + \gamma_0^2 B(x) + \gamma_1 B(-x)}{B'(-x) - \gamma_0 \{B(x) + B(-x)\}}.$$

For the case $\gamma_0 = \gamma_2 = 0$, we have $B(x) = C_1 \sinh \sqrt{2\gamma_1} x + C_2$,

hence

$$K(x, y) = A(x)B(y) + C(x) = -\sqrt{\frac{\gamma_1}{2}} \frac{\sinh \sqrt{2\gamma_1} y + \sinh \sqrt{2\gamma_1} x}{\cosh \sqrt{2\gamma_1} x}.$$

We also can show, in this case, $\gamma_3 = \frac{\gamma_1^2}{3!}$, etc., ... ; we shall discuss the converse in Example II.

In the following we shall show some examples which can be solved by Theorems 1-3.

Example I. If $R(t) = be^{-ct}$ (b, c: constants)

then

$$r_+(\alpha) = \int_0^\infty R(t)e^{i\alpha t} dt = \frac{-b}{i\alpha - c}, \quad \text{Im}(\alpha) \gg 0$$

$$r_-(\alpha) = \int_{-\infty}^0 R(t)e^{i\alpha t} dt = \frac{b}{i\alpha - c}, \quad \text{Im}(\alpha) \ll 0.$$

Hence

$$\hat{R}_+(\alpha) = \frac{r_+(\alpha) [r_(-\alpha) - 1]}{1 + r_+(\alpha) r_(-\alpha)} = \frac{(i\alpha + c + b)b}{\alpha^2 + c^2 - b^2}, \quad \text{Im}(\alpha) \gg 0$$

$$\hat{R}_-(\alpha) = \frac{-r_-(\alpha) [r_+(-\alpha) + 1]}{1 + r_+(-\alpha) r_-(\alpha)} = \frac{(i\alpha + c + b)b}{\alpha^2 + c^2 - b^2}, \quad \text{Im}(\alpha) \ll 0.$$

It follows that

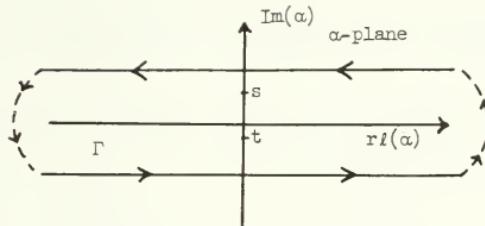
$$B(y) = \frac{1}{2\pi} \int_{is-\infty}^{is+\infty} \hat{B}_+(\alpha) e^{-i\alpha y} d\alpha + \frac{1}{2\pi} \int_{it-\infty}^{it+\infty} \hat{B}_-(\alpha) e^{-i\alpha y} d\alpha$$

($s \gg 0, t \ll 0$)

$$= \frac{1}{2\pi} \int_{\Gamma} \frac{(i\alpha+c+b)b}{\alpha^2+c^2-b^2} e^{-i\alpha y} d\alpha \quad \begin{cases} d^2 = c^2 - b^2 > 0 \\ \text{similarly for other cases.} \end{cases}$$

$$= \frac{-1}{4\pi i} \int_{\Gamma} \left\{ \frac{b}{d} \frac{d-c-b}{(\alpha-id)} + \frac{b}{d} \frac{d+c+b}{(\alpha+id)} \right\} e^{-i\alpha y} d\alpha$$

$$= \frac{-b}{2d} \left\{ (d-c-b) e^{dy} + (d+c+b) e^{-dy} \right\}$$



Γ is the contour consisting of two parallel lines as shown above.

Trying the form $K(x,y) = A(x)B(y)$ for a solution of equation (1),

we obtain exactly formula (26) if we put $\gamma_0 = b$, $\gamma_1 = -bc$.

Substituting $K(x,y) = \frac{-b}{2d} A(x) \left\{ (d-c-b)e^{dy} + (d+c+b)e^{-dy} \right\}$ into

(1), and solving for $A(x)$, we obtain (26). In the above case, of course, we could assume $K(x,y) = \sum_{n=0}^{\infty} A_n(y)x^n$ and apply Theorem 1 for

computation of $A_n(y)$, step by step; nevertheless, the computation is not very simple. For example, we have

$$A_0(y) = B(y) = \frac{-b}{2d} \left\{ (d-c-b)e^{dy} + (d+c+b)e^{-dy} \right\}$$

$$R^1(x, t) = \frac{1}{x} \left\{ b \int_{-t}^x e^{-c(t+y)} A_0(y) dy + A_0(t) + be^{-c(t+x)} \right\}$$

$$= \frac{b}{x} \left\{ \int_0^x e^{-c(t+y)} A_0(y) dy + e^{-c(t+x)} - e^{-ct} \right\}$$

since $\int_{-t}^0 be^{-c(t+y)} A_0(y) dy + A_0(t) + be^{-ct} = 0$.

It follows that

$$R^1(0, t) = b \left\{ e^{-c(t+x)} A_0(x) \Big|_{x=0} + (-c)e^{-ct} \right\}$$

$$= -b(b+c) e^{-ct}.$$

Therefore

$$r_+^1(\alpha) = \int_0^\infty R^1(0, t) e^{i\alpha t} dt = \frac{b(b+c)}{i\alpha - c}, \quad \text{Im}(\alpha) \gg 0$$

$$r_-^1(\alpha) = \int_{-\infty}^0 R^1(0, t) e^{i\alpha t} dt = \frac{-b(b+c)}{i\alpha - c}, \quad \text{Im}(\alpha) \ll 0.$$

Consequently

$$A_1(y) = \frac{-i(b+c)}{2\pi} \int_{\Gamma} \frac{i\alpha+b+c}{\alpha^2+c^2-b^2} e^{-iy\alpha} d\alpha \quad \begin{bmatrix} \Gamma \text{ is the same contour} \\ \text{as in p. 20.} \end{bmatrix}$$

$$= \frac{b(b+c)}{2d} \left\{ (d-c-b)e^{dy} + (d+c+b)e^{-dy} \right\}.$$

Similarly we can compute $A_2(y)$ as follows:

$$\begin{aligned}
 A_1(y) &= \frac{b(b+c)}{2d} \left\{ (d-c-b)e^{dy} + (d+c+b)e^{-dy} \right\} = -(b+c)A_0(y) \\
 R^2(x, t) &= \frac{1}{x} \left\{ b \int_{-t}^x e^{-c(t+y)} A_1(y) dy + A_1(t) + R^1(x, t) \right\} \\
 &= \frac{1}{x} \left\{ b \int_0^x e^{-c(t+y)} A_1(y) dy + R^1(x, t) - R^1(0, t) \right\}.
 \end{aligned}$$

It follows that

$$\begin{aligned}
 R^2(0, t) &= be^{-ct} A_1(0) + \frac{\partial}{\partial x} R^1(x, t) \Big|_{x=0} \\
 &= b^2(b+c)e^{-ct} + \frac{1}{2}b(b+c)^2 e^{-ct} = \frac{1}{2}b(b+c)(3b+c)e^{-ct}.
 \end{aligned}$$

By means of the same procedure as we used in computing $A_1(y)$, we obtain

$$\begin{aligned}
 A_2(y) &= \frac{1}{4d} b(b+c)(3b+c) \left\{ (d-c-b)e^{dy} + (d+c+b)e^{-dy} \right\} \\
 &= \frac{1}{2}(b+c)(3b+c)A_0(y).
 \end{aligned}$$

Similarly, we obtain

$$\begin{aligned}
 A_3(y) &= -\frac{1}{3!} (b+c)^2 (11b+c) A_0(y) \\
 A_4(y) &= \frac{1}{4!} (b+c)^2 (57b^2 + 38bc + c^2) A_0(y) \\
 A_5(y) &= -\frac{1}{5!} (b+c)^3 (361b^2 + 118bc + c^2) A_0(y) \quad \text{and so on.}
 \end{aligned}$$

Hence

$$K(x, y) = \sum_{n=0}^{\infty} A_n(y) x^n = A_0(y) \left\{ 1 - (b+c)x + \frac{1}{2}(b+c)(3b+c)x^2 - \dots \right\}$$

from which use of the form $K(x, y) = A(x)B(y)$ is justified.

$K(x, y)$ can be obtained from Theorem 2 (or Theorem 3) as follows:

$$\text{Since } R(t) = be^{-ct} = b \sum_{n=0}^{\infty} \frac{(-1)^n c^n}{n!} t^n, \text{ we have } \gamma_n = \frac{(-1)^n c^n}{n!} b.$$

Using the general formula for $C_{m,n}$ which we computed before, we obtain:

$$C_{0,0} = -b$$

$$C_{0,1} = b(c+b)$$

$$C_{0,2} = \frac{b}{2}(-c^2 + b^2) = -\frac{1}{2}bd^2$$

$$C_{0,3} = \frac{b}{6}(c^3 + bc^2 - b^2c - b^3) = \frac{bd^2}{3!} (c+b)$$

$$C_{0,4} = \frac{-b}{24}(c^4 - 2b^2c^2 + b^4) = \frac{-bd^4}{4!}$$

$$C_{0,5} = \frac{b}{120}(c^5 + bc^4 - 2bc^3 - 2b^2c^2 + b^3c + b^4) \\ = \frac{bd^2}{5} (c+b)^3$$

$$C_{2,0} = \frac{-b}{2}(c^2 + 4bc + 3b^2) = \frac{-b}{2}(c+b)(c+3b) \\ = \frac{1}{2}(c+b)(c+3b)C_{0,0}$$

$$C_{2,1} = \frac{b}{2}(c^3 + 5bc^2 - 7b^2c + 3b^3) \\ = \frac{b}{2}(c+b)^2(c+3b) = \frac{1}{2}(c+b)(c+3b)C_{0,1}$$

$$C_{2,2} = \frac{-b}{4}(c^4 + 4bc^3 + 2b^2c^2 - 4b^3c - 3b^2) \\ = \frac{-b}{4}d^2(c+b)(c+3b) = \frac{1}{2}(c+b)(c+3b)C_{0,2}$$

$$C_{2,3} = \frac{b}{12}(c^5 + 5bc^4 + 6b^2c^3 - 2b^3c^2 - 7b^4c - 3b^5) \\ = \frac{b}{12}d^2(c+b)^2(c+3b) = \frac{1}{2}(c+b)(c+3b)C_{0,3}$$

$$C_{4,0} = \frac{-b}{24}(c^4 + 40bc^3 + 134b^2c^2 + 152b^3c + 57b^4) \\ = \frac{-b}{4!}(c+b)^2(c^2 + 38bc + 57b^2) \\ = \frac{1}{4!}(c+b)^2(c^2 + 38bc + 57b^2)C_{0,0}$$

$$C_{4,1} = \frac{b}{24}(c^5 + 41bc^4 + 174b^2c^3 + 286b^3c^2 + 209b^4c + 57b^5) \\ = \frac{b}{4!}(c+b)^3(c^2 + 38bc + 57b^2) \\ = \frac{1}{4!}(c+b)^2(c^2 + 38bc + 57b^2)C_{1,0}$$

$$c_{1,0} = b(c+b) = -(c+b)c_{0,0}$$

$$\begin{aligned} c_{1,1} &= -b(c^2 + 2bc + b^2) + -b(c+b)^2 \\ &= -(c+b)c_{0,1} \end{aligned}$$

$$\begin{aligned} c_{1,2} &= \frac{b}{2}(c^3 + bc^2 - b^2c - b^3) = \frac{bd^2}{2}(c+b) \\ &= -(c+b)c_{0,2} \end{aligned}$$

$$\begin{aligned} c_{1,3} &= \frac{-b}{6}(c^4 + 2c^3b - 2cb^3 - b^4) \\ &= \frac{-bd^2}{3!}(c+b)^2 = -(c+b)c_{0,3} \end{aligned}$$

$$\begin{aligned} c_{1,4} &= \frac{b}{24}(c^5 + bc^4 - 2b^2c^3 - 2b^3c^2 + b^4c + b^5) \\ &= \frac{bd^4}{4!}(c+b) = -(c+b)c_{0,4} \end{aligned}$$

$$\begin{aligned} c_{3,0} &= \frac{b}{6}(c^3 + 13bc^2 + 23b^2c + 11b^3) \\ &= \frac{b}{3!}(c+b)^2(c+11b) = \frac{-1}{3!}(c+b)^2(c+11b)c_{0,0} \end{aligned}$$

$$\begin{aligned} c_{3,1} &= \frac{-b}{6}(c^4 + 14bc^3 + 36b^2c^2 + 34b^3c + 11b^4) \\ &= \frac{-b}{3!}(c+b)^3(c+11b) = \frac{-1}{3!}(c+b)^2(c+11b)c_{0,1} \end{aligned}$$

$$\begin{aligned} c_{3,2} &= \frac{b}{12}(c^5 + 13bc^4 + 22b^2c^3 - 2b^3c^2 - 23b^4c - 11b^5) \\ &= \frac{bd^2}{12}(c+1)^2(c+11b) = \frac{-1}{3!}(c+b)^2(c+11b)c_{0,2} \end{aligned}$$

$$\begin{aligned} c_{5,0} &= \frac{b}{120}(c^5 + 121bc^4 + 718b^2c^3 + 1438b^3c^2 \\ &\quad + 1201b^4c + 361b^5) \end{aligned}$$

$$\begin{aligned} &= \frac{b}{5!}(c+b)^3(c^2 + 118bc + 361b^2) \\ &= \frac{-1}{5!}(c+b)^3(c^2 + 118bc + 361b^2)c_{0,0} \end{aligned}$$

These coefficients $c_{m,n}$ are exactly the coefficients of powers $x^m y^n$ in

the expansion of $K(x, y) = -b \frac{(d-c-b)e^{dy} + (d+c+b)e^{-dy}}{(d+c+b)e^{dy} + (d-c-b)e^{-dy}}$ $\begin{cases} \text{if } d^2 = c^2 - b^2 > 0, \\ \text{similarly for } d^2 \leq 0. \end{cases}$

Example II. If $R(t) = \sqrt{b} \sinh \sqrt{b}t = \frac{\sqrt{b}}{2} (e^{\sqrt{b}t} - e^{-\sqrt{b}t})$ then

$$K(x, y) = -\sqrt{\frac{b}{2}} \frac{\sinh \sqrt{2b}x + \sinh \sqrt{2b}y}{\cosh \sqrt{2b}x}.$$

To show this, applying Theorem 1, we have

$$r_+(\alpha) = \int_0^\infty R(t) e^{i\alpha t} dt = \frac{-b}{\alpha^2 + b}, \quad \text{Im}(\alpha) \gg 0$$

$$r_-(\alpha) = \int_{-\infty}^0 R(t) e^{i\alpha t} dt = \frac{b}{\alpha^2 + b}, \quad \text{Im}(\alpha) \gg 0.$$

From the formula (7), we obtain

$$\hat{K}_+(0, \alpha) = \frac{b}{\alpha^2 + 2b}, \quad \text{Im}(\alpha) \gg 0$$

$$\hat{K}_-(0, \alpha) = \frac{-b}{\alpha^2 + 2b}, \quad \text{Im}(\alpha) \ll 0.$$

Hence

$$A_0(y) = K(0, y) = \frac{-1}{2\pi} \int_{\Gamma} \frac{be^{-iy\alpha}}{\alpha^2 + 2b} d\alpha \quad \begin{bmatrix} \Gamma \text{ is the same contour} \\ \text{as in Example I.} \end{bmatrix}$$

$$= \frac{-b}{2\sqrt{2}} \frac{1}{2\pi i} \int_{\Gamma} \left\{ \frac{1}{\alpha - i\sqrt{2b}} - \frac{1}{\alpha + i\sqrt{2b}} \right\} e^{-iy\alpha} d\alpha$$

$$= -\sqrt{\frac{b}{2}} \sinh \sqrt{2b} y.$$

Similarly, we obtain

$$A_1(y) = -b$$

* For $b = 1$, see [1], p. 19-21.

$$A_2(y) = b\sqrt{\frac{b}{2}} \sinh \sqrt{2b} y$$

$$A_3(y) = \frac{2}{3} b^2$$

$$A_4(y) = \frac{-5}{6} \sqrt{\frac{b}{2}} b^2 \sinh \sqrt{2b} y$$

$$A_5(y) = \frac{-8}{15} b^3 \quad \text{etc.}$$

Consequently,

$$K(x,y) = \sum_{n=0}^{\infty} A_n(y)x^n = -\sqrt{\frac{b}{2}} \sinh \sqrt{2b} y(1 - 6x + \frac{5}{2}b^2x^2 + \dots) + (-bx + \frac{2}{3}b^2x^3 - \frac{8}{15}b^3x^5 + \dots)$$

from which it is reasonable for us to assume that $K(x,y)$ has the form $-\sqrt{\frac{b}{2}} A(x) \sinh \sqrt{2b} y + C(x)$; substituting the latter into (1), we obtain $A(x) = \operatorname{sech} \sqrt{2b} x$, $C(x) = -\sqrt{\frac{b}{2}} \tanh \sqrt{2b} x$.

It also can be seen from Theorem 2 (or Theorem 3) that, in this case, we have $\gamma_{2n} = 0$, $\gamma_{2n+1} = \frac{b^{n+1}}{(2n+1)!}$; hence

$$\begin{array}{lllll} C_{0,2n} = 0 & C_{1,0} = -b & C_{2,0} = 0 & C_{3,0} = \frac{2}{3} b^2 & C_{4,0} = 0 \\ C_{0,0} = -b & C_{1,1} = 0 & C_{2,1} = b^2 & C_{3,1} = 0 & C_{4,1} = \frac{-5}{6} b^3 \\ C_{0,3} = -\frac{2}{3!} b^2 & C_{1,2} = 0 & C_{2,2} = 0 & C_{3,2} = 0 & C_{4,2} = 0 \\ C_{0,5} = -\frac{2^2}{5!} b^3 & C_{1,3} = 0 & C_{2,3} = \frac{2b^3}{3} & C_{3,3} = 0 & C_{4,3} = \frac{-5}{12} b^4 \end{array}$$

from which we could also guess that $K(x,y)$ is of the form

$$A(x)B(y) + C(x) \quad \text{with} \quad B(y) = -\sqrt{\frac{b}{2}} \sinh \sqrt{2b} y, \quad C(x) = -\sqrt{\frac{b}{2}} \tanh \sqrt{2b} x.$$

Substituting these into (1), we obtain $A(x) = \operatorname{sech} \sqrt{2b} x$.

From the above example, it is known that Theorem 2 is useful in determining the form of $K(x,y)$ for a given $R(t)$; similarly, for a given $r_+(a)$, Theorem 3 is also useful in determining the form of $K(x,y)$.

As in Example II, it is found that if $R(t)$ is an odd function (i.e.

$\gamma_{2n} = 0$), then $C_{m,n} = 0$ for $m+n$ is even; and hence $K(x,x)$ is also an odd function. Consequently, $V(x) = 2 \frac{d}{dx} K(x,x)$ is an even function.

When $R(t)$ is even we cannot conclude very much as in the following example.

Example III If $R(z) = e^{-cz^2}$, then $\gamma_{2m} = \frac{(-1)^m c^m}{m}$, $\gamma_{2m+1} = 0$.

From Theorem 2, we then obtain $C_{p,q}$ ($p+q \leq 5$) as follows:

$$C_{0,0} = -1$$

$$C_{0,1} = 1$$

$$C_{0,2} = c + \frac{1}{2}$$

$$C_{0,3} = -\frac{2}{3}c - \frac{1}{6}$$

$$C_{0,4} = -\frac{c^2}{2} - \frac{1}{4}c - \frac{1}{24}$$

$$C_{0,5} = \frac{7}{30}c^2 + \frac{1}{15}c + \frac{1}{120}$$

$$C_{2,0} = c - \frac{3}{2}$$

$$C_{2,1} = -2c + \frac{3}{2}$$

$$C_{2,2} = -3c^2 + \frac{1}{2}c + \frac{3}{4}$$

$$C_{2,3} = \frac{7}{3}c^2 - \frac{2}{3}c - \frac{1}{4}$$

$$C_{1,0} = 1$$

$$C_{1,1} = 2c-1$$

$$C_{1,2} = -\frac{1}{2}$$

$$C_{1,3} = -2c^2 + \frac{1}{3}c + \frac{1}{6}$$

$$C_{1,4} = \frac{-1}{6}c^2 + \frac{1}{6}c + \frac{1}{24}$$

$$C_{3,1} = -\frac{8}{3}c + \frac{11}{6}$$

$$C_{3,1} = -2c^2 + \frac{13}{3}c - \frac{11}{6}$$

$$C_{3,2} = \frac{7}{3}c^2 + \frac{1}{3}c - \frac{11}{12}$$

$$C_{4,0} = -\frac{c^2}{2} + \frac{53}{12}c - \frac{19}{8}$$

$$C_{4,1} = \frac{23}{6}c^2 - \frac{19}{3}c + \frac{19}{8}$$

*See Appendix

$C_{5,0} = \frac{127}{30}c^2 - \frac{223}{30}c + \frac{361}{120}$ and so on.

And

$$K(x, x) = -1 + 2x + 2(2c - 1)x^2 + \frac{16}{3} (c + \frac{1}{2})x^3 + (-8c^2 + \frac{28}{3}c - \frac{10}{3})x^4 + (\frac{64}{5}c^2 - \frac{208}{15}c + \frac{64}{15})x^5 + \dots$$

Example IV

$$\text{If } R(z) = ze^{-cz^2}, \text{ then } \gamma_{2m} = 0, \gamma_{2m+1} = \frac{(-1)^m c^m}{m!} \quad .$$

And we obtain $c_{m,n} = 0$ if $m+n$ is even.

$$c_{0,1} = -1, \quad c_{1,0} = -1$$

$$c_{2,1} = 3c + \frac{3}{2}$$

$$c_{0,3} = c - \frac{1}{6}, \quad c_{1,2} = 3c + \frac{1}{2}$$

$$c_{2,3} = -5c^2 - c + \frac{1}{4}$$

$$c_{0,5} = \frac{-c^2}{2} + \frac{c}{10} - \frac{1}{120}, \quad c_{1,4} = -\frac{5c^2}{2} - \frac{c}{2} - \frac{1}{24}, \quad c_{3,0} = -c + \frac{5}{6}$$

$$c_{3,2} = -5c^2 - 3c - \frac{5}{12}, \quad c_{4,1} = -\frac{5c^2}{2} - \frac{11}{2}c - \frac{41}{24}, \quad c_{5,0} = -\frac{c^2}{2} - \frac{29}{10}c - \frac{121}{120};$$

$$\text{hence } K(x, x) = -2x + (6c + \frac{16}{6})x^3 - (16c^2 + \frac{64}{5}c + \frac{44}{15})x^5 + \dots$$

$$\text{and } V(x) = 2 \frac{d}{dx} K(x, x) = -4 + 6(6c + \frac{16}{6})x^2 - 10(16c^2 + \frac{64}{5}c + \frac{44}{15})x^4 + \dots$$

APPENDIX

(i) holds since

$$\begin{aligned} \int_0^\infty \int_0^w R(w-u) K(x, x-u) du e^{i\alpha w} dw &= \int_0^\infty K(x, x-u) e^{i\alpha u} du \int_u^\infty R(w-u) e^{i\alpha(w-u)} dw \\ &= r_+(\alpha) \int_0^\infty K(x, x-u) e^{i\alpha u} du = r_+(\alpha) \int_{-\infty}^x K(x, y) e^{-i\alpha y} dy \cdot e^{i\alpha x} \\ \int_0^\infty K(x, w-x) e^{i\alpha w} dw &= \int_x^\infty K(x, y) e^{i\alpha y} dy \cdot e^{i\alpha x}. \end{aligned}$$

Similarly, (ii) holds since

$$\begin{aligned} \int_{-\infty}^0 \int_0^w R(w-u) K(x, x-u) du e^{i\alpha w} dw &= - \int_{-\infty}^0 K(x, x-u) e^{i\alpha u} du - \int_{-\infty}^u R(w-u) e^{i\alpha(w-u)} dw \\ &= -r_-(\alpha) \int_{-\infty}^0 K(x, x-u) e^{i\alpha u} du = -r_-(\alpha) \int_x^{-\infty} K(x, y) e^{-i\alpha y} dy \cdot e^{i\alpha x} \\ \int_{-\infty}^0 K(x, w-x) e^{i\alpha w} dw &= \int_{-\infty}^{-x} K(x, y) e^{i\alpha y} dy \cdot e^{i\alpha x}. \end{aligned}$$

The identity (*) can be derived as follows:

$$(1-x)^b = \sum_{s=0}^b (-1)^s \binom{b}{s} x^s, \quad (1-x)^{-a} = \sum_{r=0}^{\infty} \binom{a+r-1}{a-1} x^r$$

and

$$(1-x)^{b-a} = \sum_{t=0}^{b-a} (-1)^t \binom{b-a}{t} x^t \quad (\text{we assume } b \geq a, b, a \text{ are natural numbers})$$

It follows that by equating the t^{th} power of $(1-x)^{b-a}$ and the product of $(1-x)^b$ and $(1-x)^{-a}$, we obtain

$$(-1)^t \binom{b-a}{t} = \sum_{s+r=t} (-1)^s \binom{b}{s} \binom{a+r-1}{a-1} = \sum_{s=0}^t (-1)^s \binom{b}{s} \binom{a-1+t-s}{a-1},$$

putting $a = k+1$, $b = l+1$, $t = l-k$, we obtain

$$(**) \quad (-1)^{l-k} = \sum_{s=0}^{l-k} (-1)^s \binom{l+1}{s} \binom{l-s}{k} \quad .$$

And since

$$\binom{l+1}{s} \binom{l-s}{k} = \frac{(l+1)!}{(l+1-s)! s!} \frac{(l-s)!}{(l-s-k)! k!} = \frac{l+1}{l-s+1} \frac{l!}{k! (l-k)!} \frac{(l-k)!}{(l-s-k)! s!}$$

$$= (l+1) \binom{l}{k} \frac{\binom{l-k}{s}}{l-s+1} \quad ,$$

by substituting the righthand side of the above into $(**)$ and dividing by $(-1)^{l-k}$, we obtain the identity $(*)$ easily.

To prove $C_{m,n} = 0$ for $m+n = \text{even}$, if $\gamma_{2m} = 0$:

From (18), it is easily shown that $C_{m,n}$ has at most the following terms:

$$\gamma_{m+n}, \gamma_{m+n-1}\gamma_0, \gamma_{m+n-2}\gamma_1, \gamma_{m+n-2}\gamma_0^2, \dots$$

These are isomorphic to all partitions of $m+n+1$ into natural numbers such that l corresponds to γ_{l-1} ; e.g., $(m+n+1) \leftrightarrow \gamma_{m+n}; (m+n-1, 1, 1) \leftrightarrow \gamma_{m+n-2}\gamma_0^2$ etc. Therefore, if $m+n$ is even, then $m+n+1$ is odd; hence all their partitions are of the form (p_1, p_2, \dots, p_i) such that one of p_i ($1 \leq i \leq m+n+1$) is odd. Otherwise the sum $\sum p_i$ is not an odd number since sums of even numbers are even. It follows that $C_{m,n} = 0$ (if $m+n = \text{even}$) since each of their terms has at least one factor $\gamma_{p_i-1} = 0$ which corresponds to odd p_i .

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